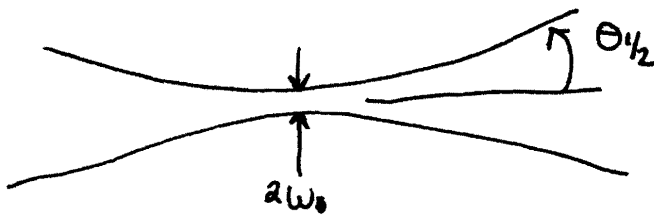


$$\theta_d \sim \frac{\lambda}{D}$$

DIVERGENCE ANGLE - MOST OF LIGHT WITHIN CONE

FOR GAUSSIAN LASER BEAM



$w_0 \equiv$ INITIAL BEAM RADIUS

$$\theta_{1/2} = \frac{\lambda}{\pi w_0} \quad \text{- FROM GAUSSIAN BEAM PROPAGATION DERIVATION}$$

ACTUAL LASERS CHARACTERIZED BY $\theta_c \sim \frac{\lambda}{D_c}$

$D_c \leftarrow$ TRANSVERSE COHERENCE LENGTH

INCOHERENT \equiv SMALL D_c , RAPID SPREAD $D_c \ll D$

COHERENT $D_c = D$ - BEST YOU CAN DO - "DIFFRACTION LIMITED"

DEFINE $M^2 = \frac{\text{DIVERGENCE OF ACTUAL LASER}}{\text{DIVERGENCE OF IDEAL LASER}} = \frac{\theta_c}{\theta_d} = \frac{D}{D_c}$

M^2 CLOSER TO 1.0 MEANS CLOSER TO IDEAL LASER.

GAUSSIAN BEAMS

MAXWELL'S EQUATIONS CAN BE RE-ARRANGED TO YIELD THE SO-CALLED "WAVE EQUATION"

$$\nabla^2 \vec{E}(\vec{r}, t) = \frac{1}{v^2} \frac{\partial^2 \vec{E}}{\partial t^2} \quad \text{where } v = \frac{c}{n} \quad (1)$$

TRIAL SOLUTION OF FORM $E(\vec{r}, t) = U(x, y, z) e^{i(kz - \omega t)}$ (2)

\uparrow SOME FUNCTION WHICH DETERMINES TRANSVERSE + SPATIAL PROFILE
 \uparrow PLANE WAVE

SUB OUR GUESS INTO WAVE EQUATION (1)

$$e^{i(kz - \omega t)} \left[\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} + \frac{\partial^2 U}{\partial z^2} + 2ik \frac{\partial U}{\partial z} - \left(k^2 - \frac{\omega^2}{c^2} \right) U \right] = 0$$

\uparrow SINCE ATT. VS Z IS SMALL
 \uparrow $\omega = kc$ SO TERM IS 0

$$\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} + 2ik \frac{\partial U}{\partial z} = 0 \quad (3)$$

MAKE EDUCATED GUESS FOR U BASED ON EXPECTED TRANSVERSE BEHAVIOUR

$$\tilde{U}(x, y, z) = E_0 e^{i(p(z) + k[x^2 + y^2]/2q(z))}$$

WHERE $p(z)$ AND $q(z)$ ARE GENERAL, UNKNOWN FUNCTIONS.

PLUGGING INTO (3) AND REARRANGING

$$\left[\underbrace{\left(\frac{2ik}{q} - 2k \frac{\partial p}{\partial z} \right)}_{\uparrow} + \underbrace{\left(\frac{k^2}{q^2} \frac{\partial q}{\partial z} - \frac{k^2}{q^2} \right)}_{\uparrow} (x^2 + y^2) \right] \tilde{U} = 0$$

CAN TREAT THESE COEFFICIENTS SEPERATELY

SO $\frac{\partial p}{\partial z} = \frac{i}{q}$ AND $\frac{\partial q}{\partial z} = 1$

MAKE A GOOD GUESS THAT $\frac{1}{q} = \frac{1}{R} + i \frac{\lambda}{\pi w^2}$ (1) (2) (3)

WHERE R IS RADIUS OF CURVATURE

w IS "SPOT SIZE" IS BEAMS HALF-WIDTH WHERE IRRADIANCE HAS DROPPED TO $1/e^2$ OF ON-AXIS VALUE

PLUGGING (1) INTO (2) GIVES

$$\tilde{U}(x, y, z) = E_0 e^{i p(z)} e^{i k(x^2 + y^2) / 2R(z)} e^{-(x^2 + y^2) / w^2}$$

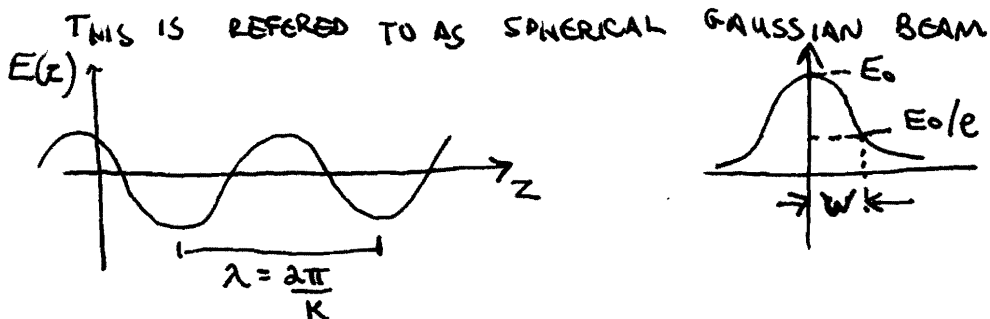
SO FROM (2)

$$E(x, y, z, t) = E_0 e^{i k(x^2 + y^2) / 2R(z)} e^{-(x^2 + y^2) / w^2(z)} e^{i [kz + p(z) - \omega t]}$$

DESCRIBES THE
WAVEFRONT

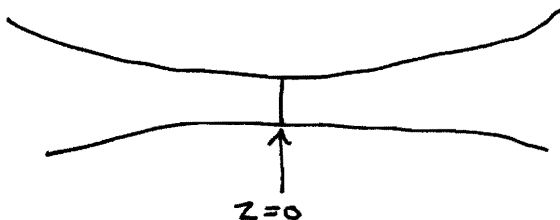
DESCRIBES
TRANSVERSE
PLANE

Z-DEPENDENT
PHASE INFORMATION
 $p(z)$ YET TO BE DETERMINED



HOW CAN WE USE THIS SOLUTION?

SUPPOSE AT SOME POINT, $z=0$ WE HAVE A PLANAR WAVEFRONT



THEN $R \rightarrow \infty$

SO FROM (4) $\frac{1}{\tilde{q}_0} = \frac{i\lambda}{\pi W_0^2}$ OR $\tilde{q}_0 = -\frac{i\pi W_0^2}{\lambda}$

NOW, SINCE $\frac{d\tilde{q}}{dz} = 1$ THEN $\tilde{q}(z) = \tilde{q}(0) + z$ (SINCE $q(0) = q_0$)

$$\text{SO } \tilde{q}(z) = z - i \frac{\pi W_0^2}{\lambda}$$

WRITING IN RECIPROCAL FORM, AND RATIONALIZING DENOMINATOR

$$\frac{1}{\tilde{q}(z)} = \frac{z}{\left(z^2 + \frac{\pi^2 W_0^4}{\lambda^2}\right)} + i \frac{\frac{\pi W_0^2}{\lambda}}{\left(z^2 + \frac{\pi^2 W_0^4}{\lambda^2}\right)}$$

COMPARE THIS TO OUR INITIAL GUESS FOR $\frac{1}{\tilde{q}(z)}$

$$\frac{1}{\tilde{q}(z)} = \frac{1}{R} + i \frac{\lambda}{\pi W^2}$$

IT CAN BE SHOWN

$$R(z) = z \left[1 + \left(\frac{\pi W_0^2}{\lambda z} \right)^2 \right] \quad (5)$$

$$W(z) = W_0^2 \left[1 + \left(\frac{\lambda z}{\pi W_0^2} \right)^2 \right] \quad (6)$$

WHERE $R(z)$ AND $W(z)$ ARE ACTUAL PHYSICAL PARAMETERS

NOTE: IF $R(z) = z$ ONLY, THEN WE WOULD HAVE SPHERICAL WAVEFRONTS CENTRED AT $z=0$.

HOWEVER IF $z \gg \frac{\pi W_0^2}{\lambda}$ THEN $R(z) = z$ AND WE DO HAVE \uparrow . THIS IS CALLED "FAR FIELD"

AND ALSO $W(z) = \frac{\lambda z}{\pi W_0}$ (LINEAR WITH z), $\theta_{FF} = \frac{\lambda}{\pi W_0}$